

# Krylov Winding and Emergent Coherence in Operator Growth Dynamics

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Thomas Scaffidi,  
University of California, Irvine

# Operator Dynamics

- Dynamics of an operator in Heisenberg picture  $O(t) = e^{iHt} O e^{-iHt}$
- 2- and 4-point functions (including OTOCs) capture only a small fraction of the information
- We want to study the *operator wave function* in a suitable basis

$$|O(t)\rangle = \sum_P c_P(t) |P\rangle$$

Eg. Pauli basis, Krylov basis, ...

# Operator Space: Size basis

- Pauli (or size) basis  $\sigma_1^\alpha \otimes \sigma_2^\alpha \otimes \dots \otimes \sigma_N^\alpha \equiv P$   $\sigma^\alpha \in \{\mathbb{1}, X, Y, Z\}$
- Size  $|P|$  = number of **non-identities**  $|Y_2 \otimes Z_3| = 2, |X_1 \otimes Z_2 \otimes Y_3| = 3$

$$O(t) = X_1 + 2.1Y_2 \otimes Z_3 - 0.8Z_2 \otimes Y_3$$

$$O(t) = e^{iHt} O e^{-iHt}$$



$$|O(t)\rangle = |X_1\rangle + 2.1|Y_2 \otimes Z_3\rangle - 0.8|Z_2 \otimes Y_3\rangle$$

$$|O(t)\rangle = e^{i\mathcal{L}t} |O\rangle$$

# Operator space: Krylov basis

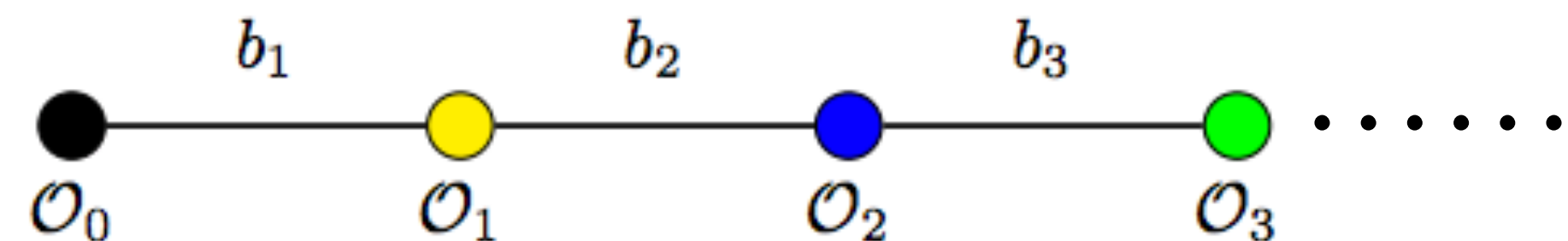
- $|O(t)\rangle = e^{i\mathcal{L}t}|O\rangle = |O\rangle + (it)\mathcal{L}|O\rangle + \frac{(it)^2}{2!}\mathcal{L}^2|O\rangle + \dots$
- $|O(t)\rangle$  is an element in the space spanned by the basis  $\{|O\rangle, \mathcal{L}|O\rangle, \mathcal{L}^2|O\rangle, \dots\}$
- Perform Gram-Schmidt on this set to obtain the **Krylov basis**  $\{|O_0\rangle, |O_1\rangle, |O_2\rangle, \dots\}$
- The Liouvillian is tridiagonal in the Krylov basis  $\{|O_n\rangle\}_{n=0}^{\infty}$

$$(O_1 | O_2) = \text{Tr}[O_1^\dagger O_2]$$

$$(O_n | \mathcal{L} | O_m) = \begin{pmatrix} 0 & b_1 & 0 & 0 & \dots \\ b_1 & 0 & b_2 & 0 & \dots \\ 0 & b_2 & 0 & b_3 & \dots \\ 0 & 0 & b_3 & 0 & \ddots \\ \vdots & \vdots & \vdots & \ddots & \ddots \end{pmatrix}$$

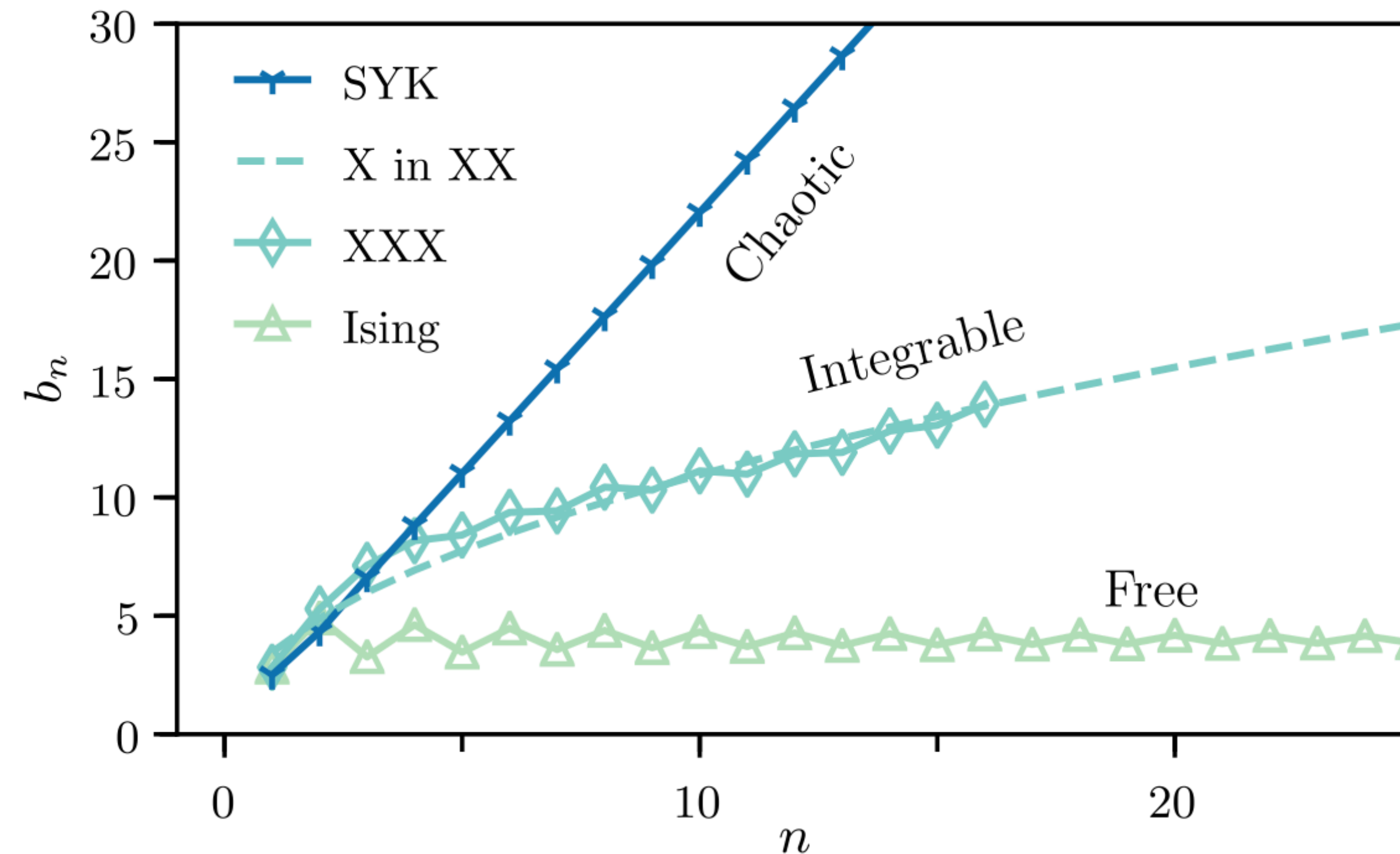
$$|O(t)\rangle = \sum_n \varphi_n(t) |O_n\rangle$$

$$\partial_t \varphi_n(t) = b_n \varphi_{n-1}(t) - b_{n+1} \varphi_{n+1}(t)$$



# Operator Growth Hypothesis

Asymptotics ( $n \gg 1$ )



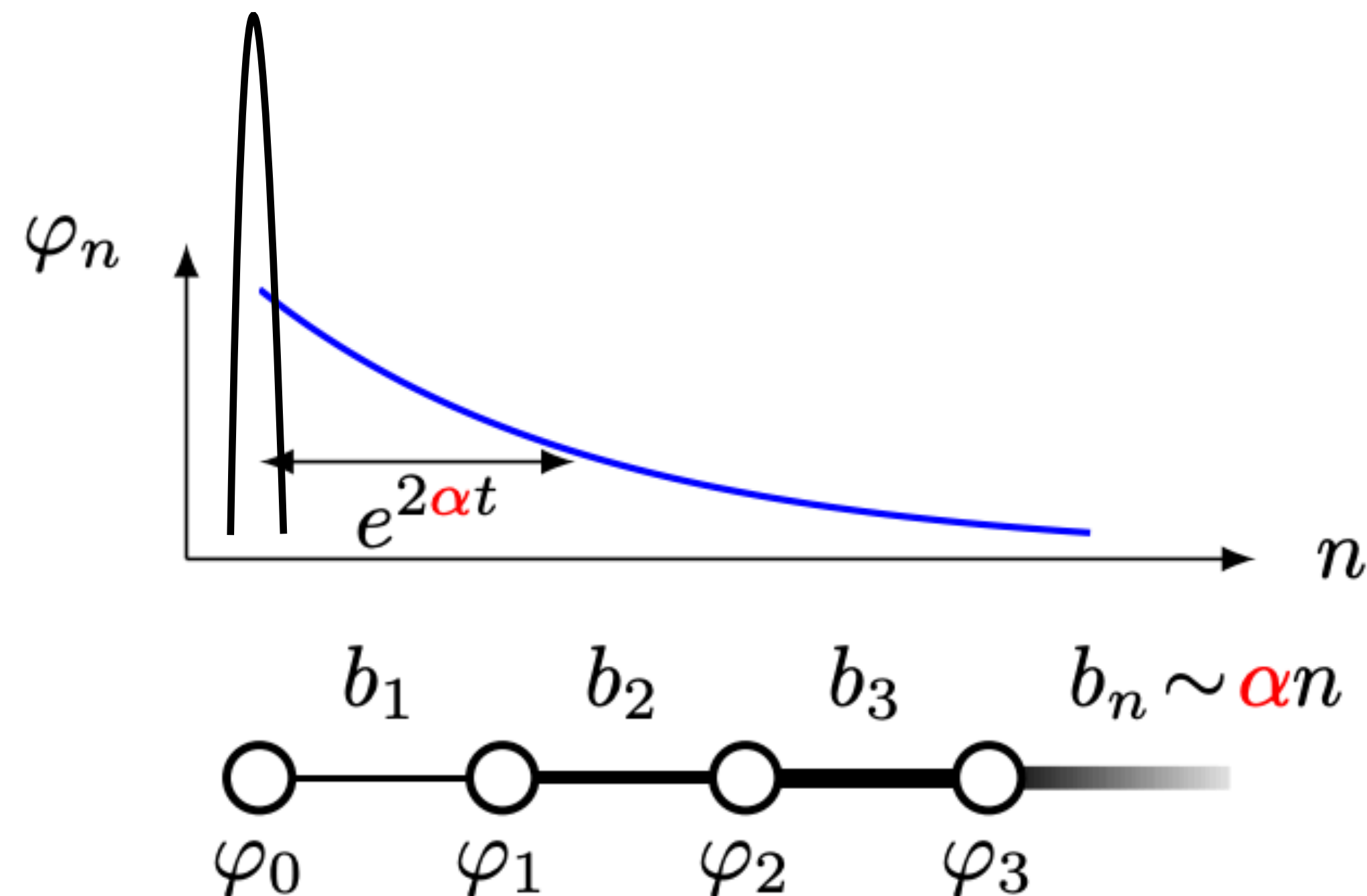
**Chaotic models**  $b_n \sim n$

**Integrable models**  $b_n \sim \sqrt{n}$

**Free models**  $b_n \sim O(1)$

# Operator Growth Hypothesis

$t = 0$



- Escape to infinity exponentially fast

$$b_n \xrightarrow{n \gg 1} \alpha n + \gamma \implies \varphi_n(t) \sim e^{-\frac{n}{e^{2\alpha t}}}$$

- Krylov complexity = average position on the chain

$$n(t) = \sum_n n \varphi_n(t)^2 \sim e^{2\alpha t}$$

All to all k-local models

What happens to the operator  
dynamics at finite temperature

$$T < \infty?$$

# Krylov Winding

- We want to study  $\rho_\beta^{1/2} O(t)$  where  $\rho_\beta = e^{-\beta H} / Z$
- Provides natural route to study linear response  $C(t) = \text{Tr}[\rho_\beta O(t) O] = (\rho_\beta^{1/2} O | \rho_\beta^{1/2} O(t))$

$$|\rho_\beta^{1/2} O(t)\rangle = e^{i\mathcal{L}(t+i\beta/4)} |\rho_\beta^{1/4} O \rho_\beta^{1/4}\rangle$$

- Generate Krylov basis with the seed operator  $\rho_\beta^{1/4} O \rho_\beta^{1/4}$

$$|\rho_\beta^{1/2} O(t)\rangle = \sum_n \varphi_n(t + i\beta/4) |O_n\rangle$$

- The Krylov wave function becomes **complex** and acquires a **phase**

**non-Hermitian operator  $\iff$  complex wave function**

# Krylov Winding

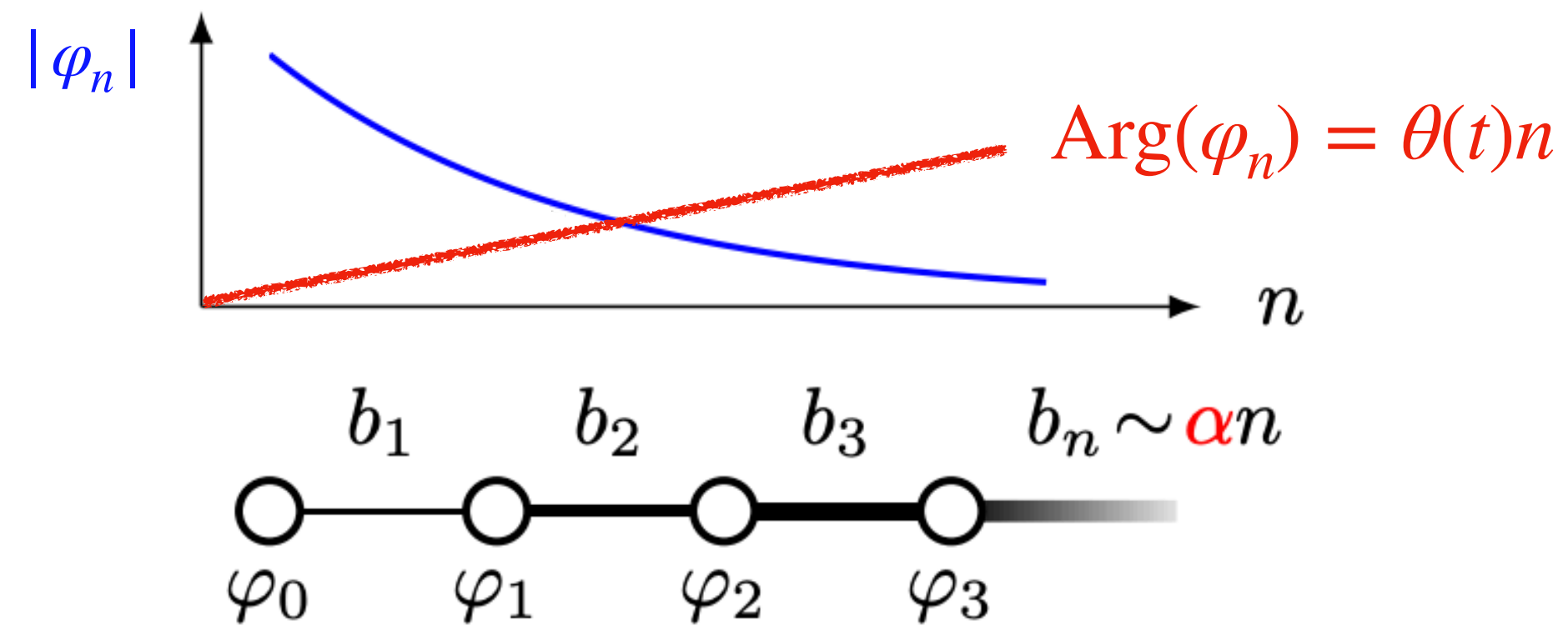
- Assuming perfectly linear  $b_n = \alpha n$ , we can use the exact solution of Krylov dynamics:

$$\varphi_n(t) = \frac{\tanh[\alpha(t + i\beta/4)]^n}{\cosh[\alpha(t + i\beta/4)]} \equiv |\varphi_n| \exp(i\theta(t)n)$$

$$\theta(t) = \tan^{-1} \left( \frac{\sin(\alpha\beta/2)}{\sinh(2\alpha t)} \right)$$

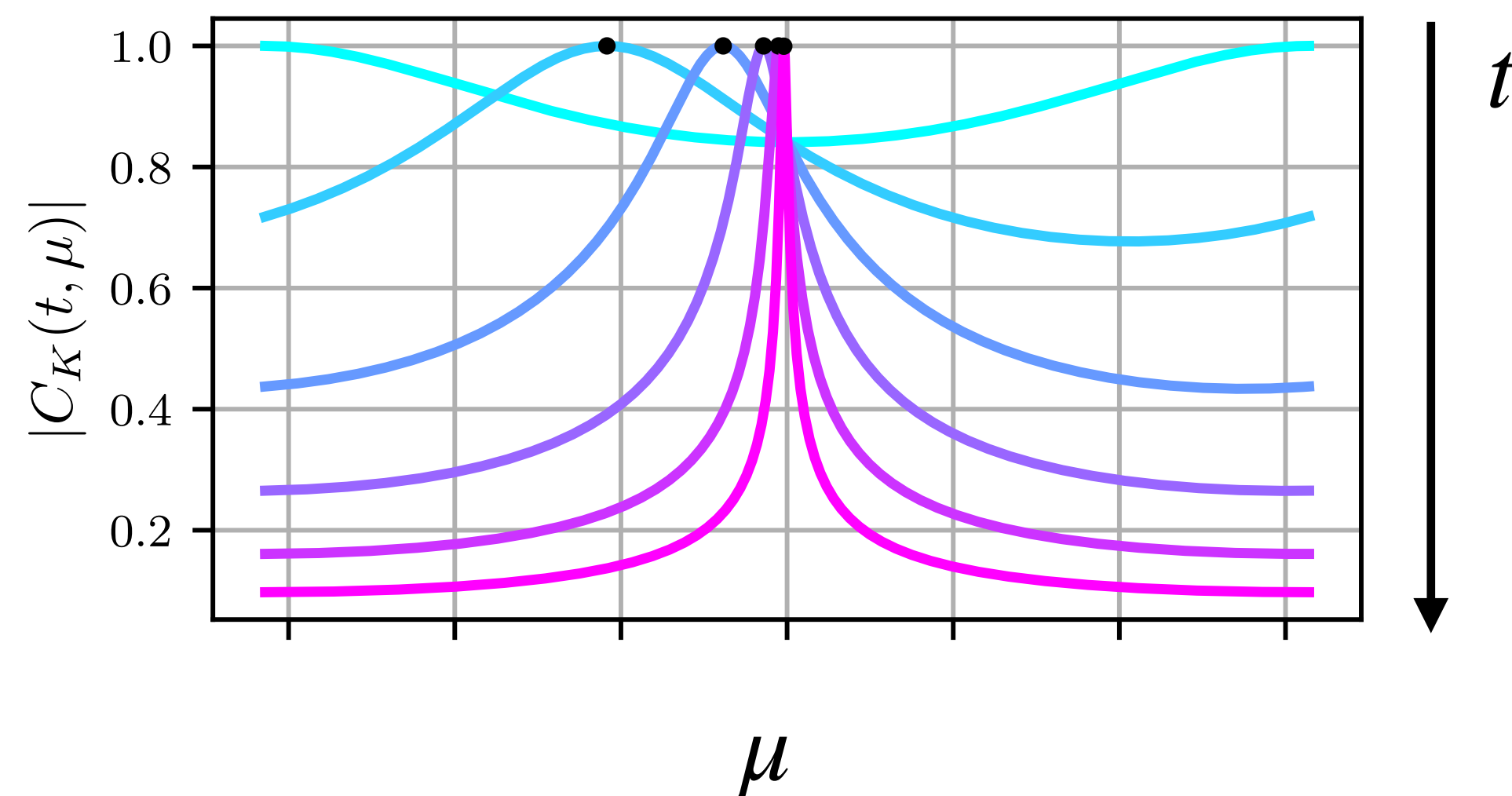
**The Krylov wave function acquires a phase that depends linearly on the Krylov index  $n$**

# Krylov Winding



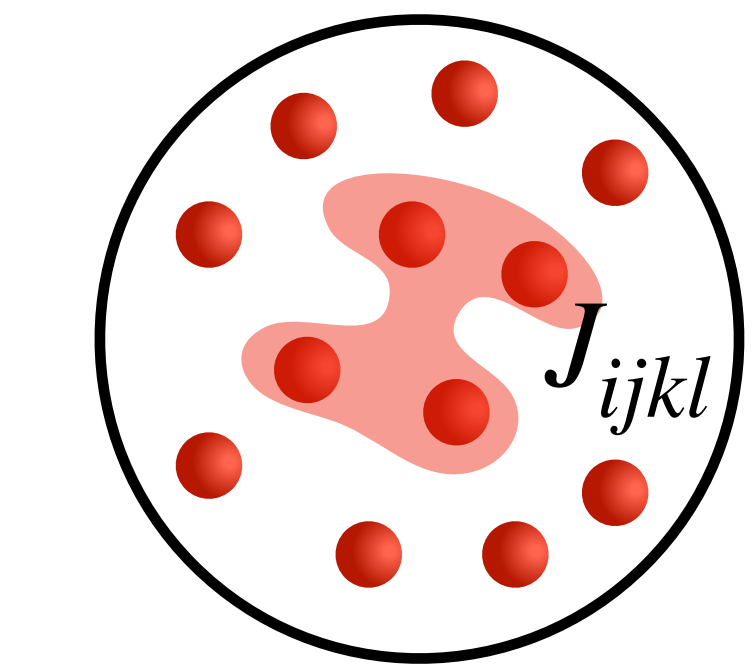
- Define *Momentum-space* Krylov wave function

$$C_K(t, \mu) = \sum_n \varphi_n^2(t) e^{i\mu n}$$



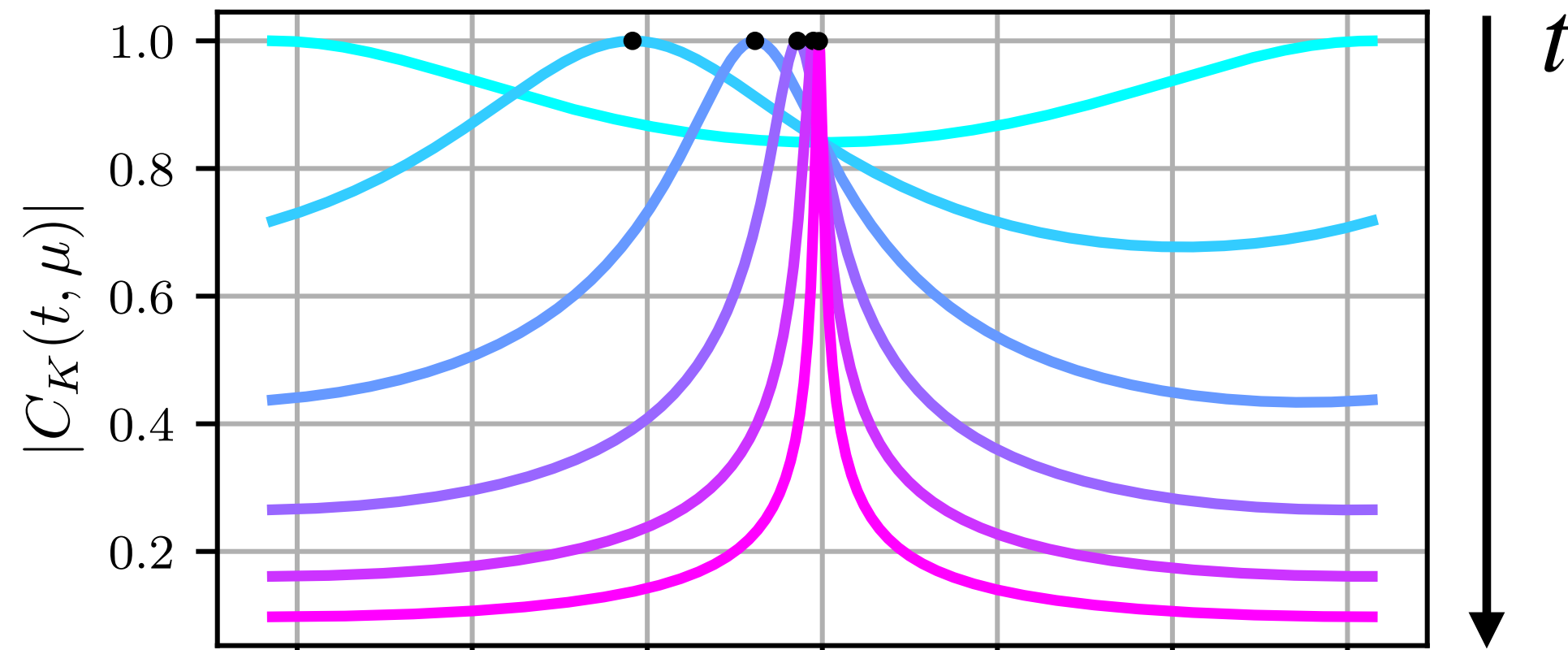
- Krylov winding in  $n$  space  $\varphi_n \sim |\varphi_n| e^{i\theta n} \Leftrightarrow$  Lorentzian peak in momentum space at  $\mu_K = -2\theta$

# Krylov Winding: Evidence



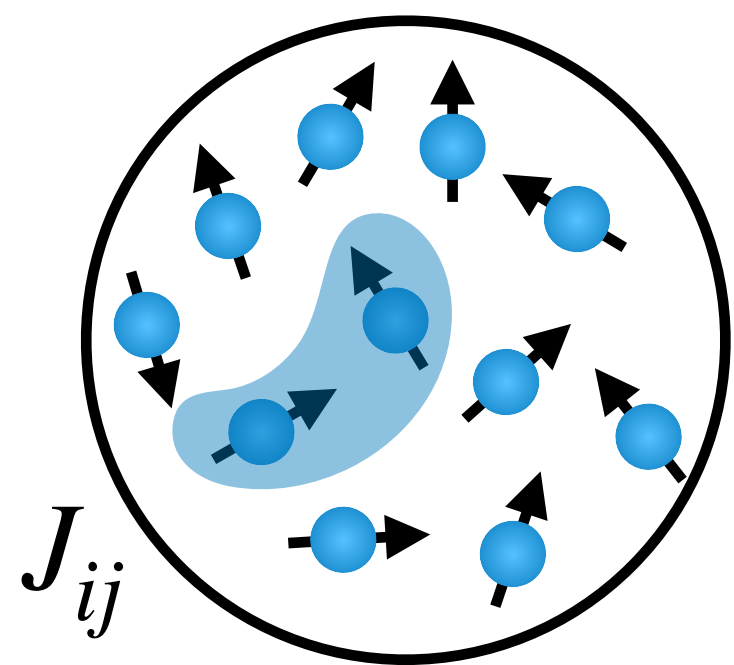
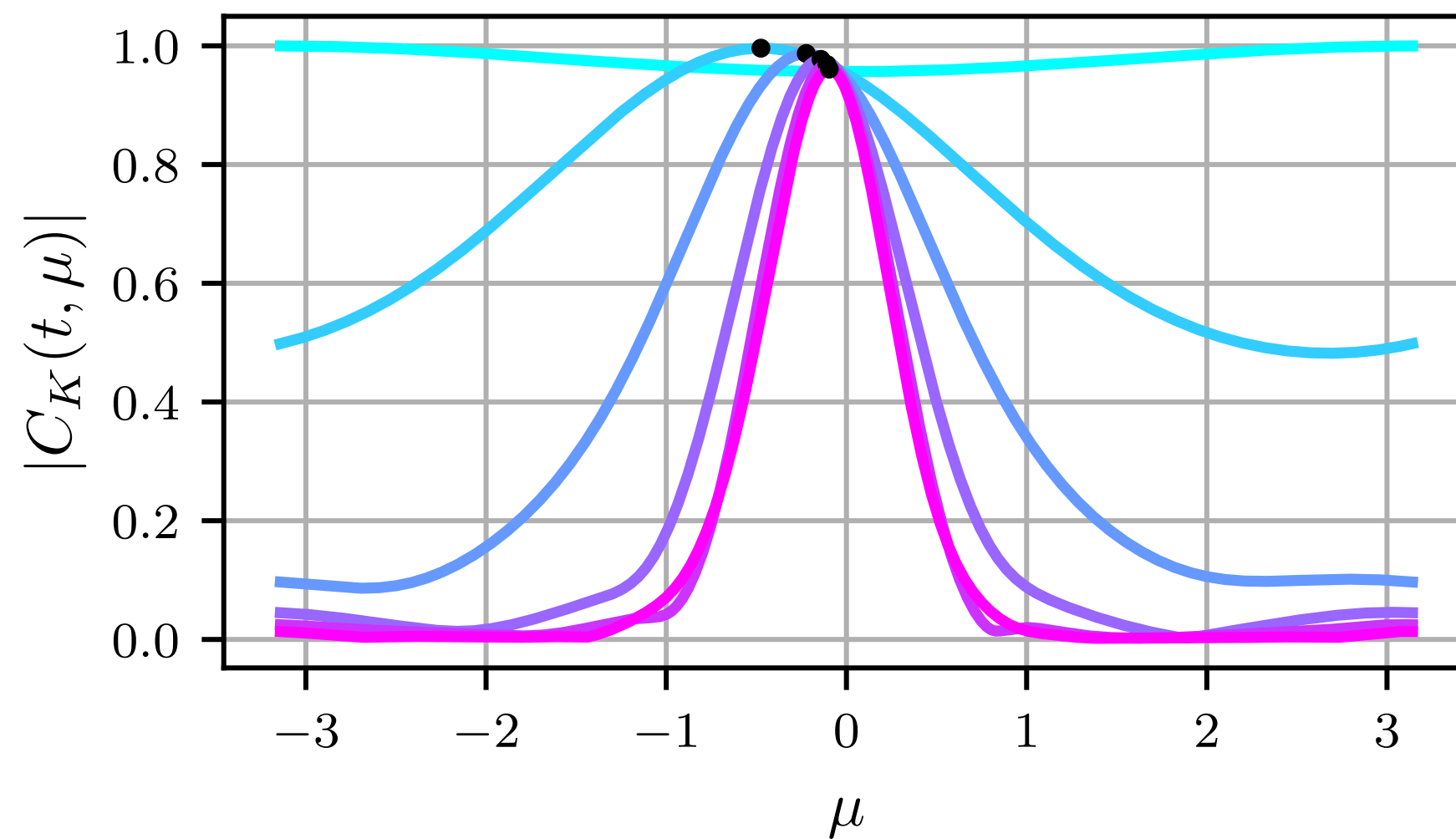
Large- $q$  SYK model

(a) Analytically solvable model

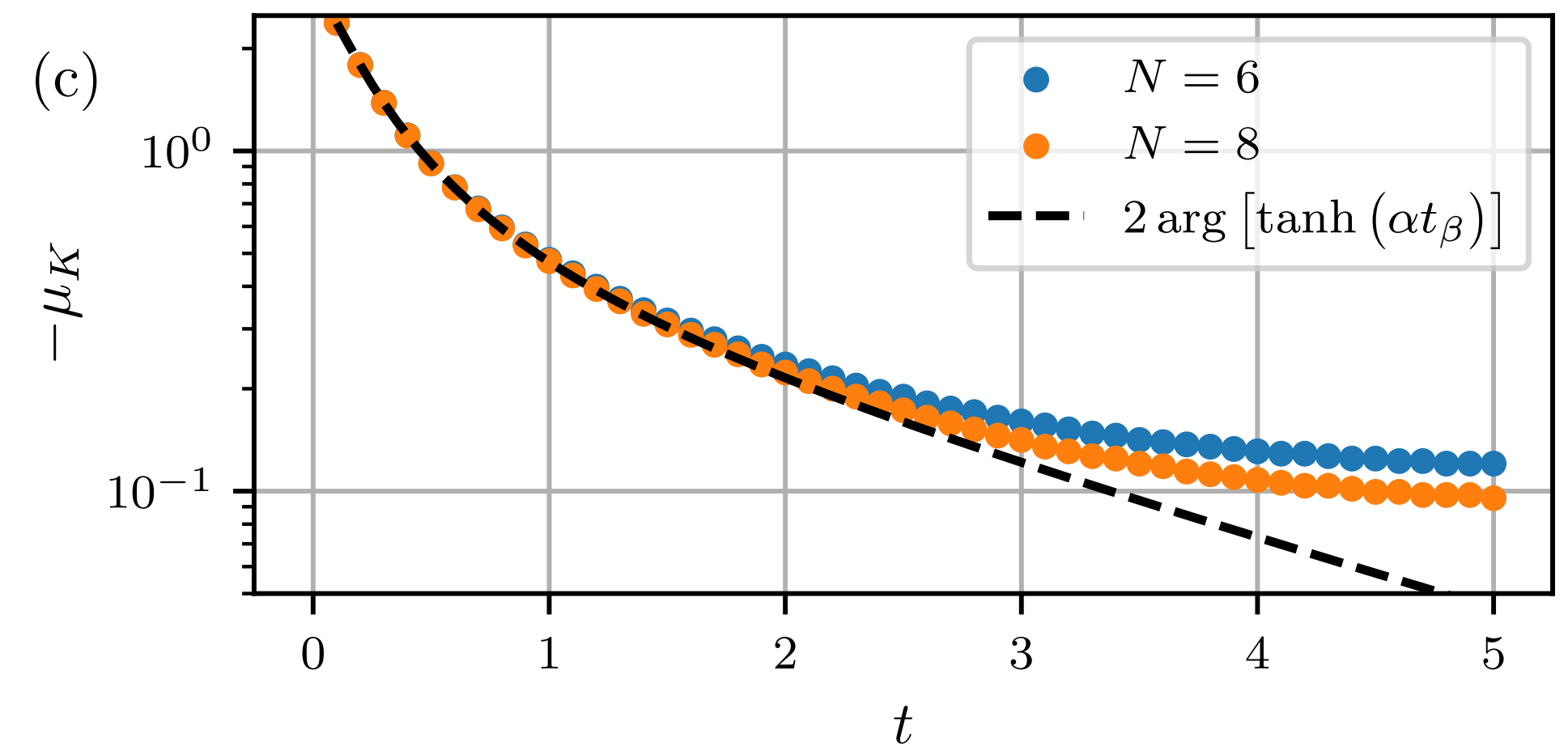


$$\mu_K(t) \sim -e^{-2\alpha t} \sin(\alpha\beta/2)$$

(b) Spin model



All-to-all Spin model



# Size Winding

Express the operator in the Pauli basis:  $|\rho^{1/2}O(t)\rangle = \sum_P c_P(t) |P\rangle$

- Because this operator is not hermitian, the  $c_P$  are in general complex
- Remarkably, certain models exhibit **size winding**:  $c_P(t) = |c_P(t)| e^{i\theta(t)|P|}$ 
  - **(1) Phase alignment: All Paulis of the same size have the same phase**
  - **(2) Phase linearity: The phase depends linearly on the Pauli size**
- Resource for “many-body teleportation”

# Size Winding: Phase alignment

**“Size-resolved Krylov overlap matrix”:**

$$M_{nm}(l) \equiv (O_n | \hat{P}_l | O_m)$$

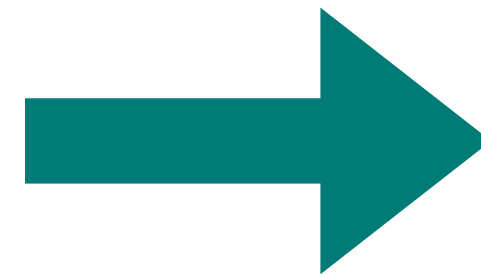
[Chen, Mu, Wang, Zhang PRL 2025]

with  $\hat{P}_l$  the projector onto the size- $l$  sector

**Rank 1**

$$M_{nm}(l) = \psi_n(l)\psi_m(l)$$

Evidence: large- $q$  SYK + Bath model



**Perfect phase alignment**

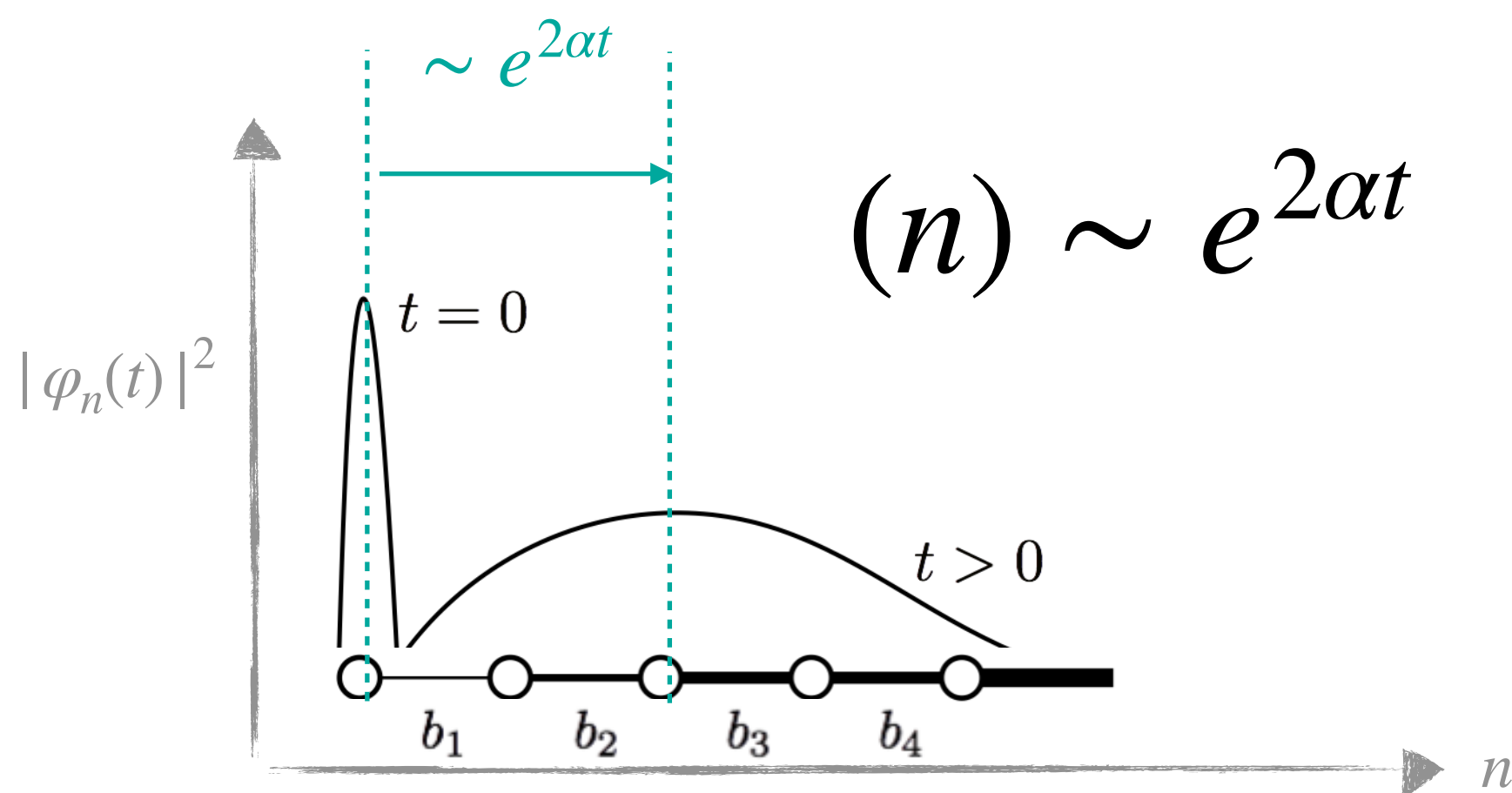
More generally, (imperfect) phase alignment is guaranteed if  $\text{rank}[M_{nm}] = \simeq 1$

(i.e. if  $M$  has a single dominant eigenvalue).

Conjecture: this is true generically for  $k$ -local models.

# Krylov vs Size growth

**Krylov wavefunction**

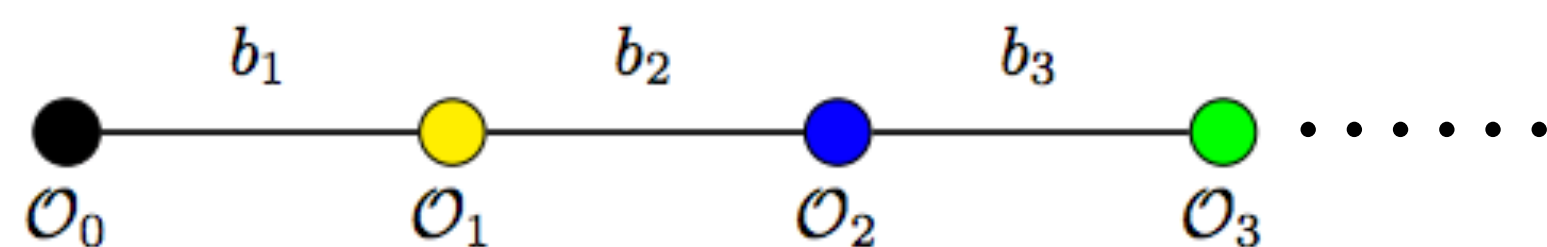
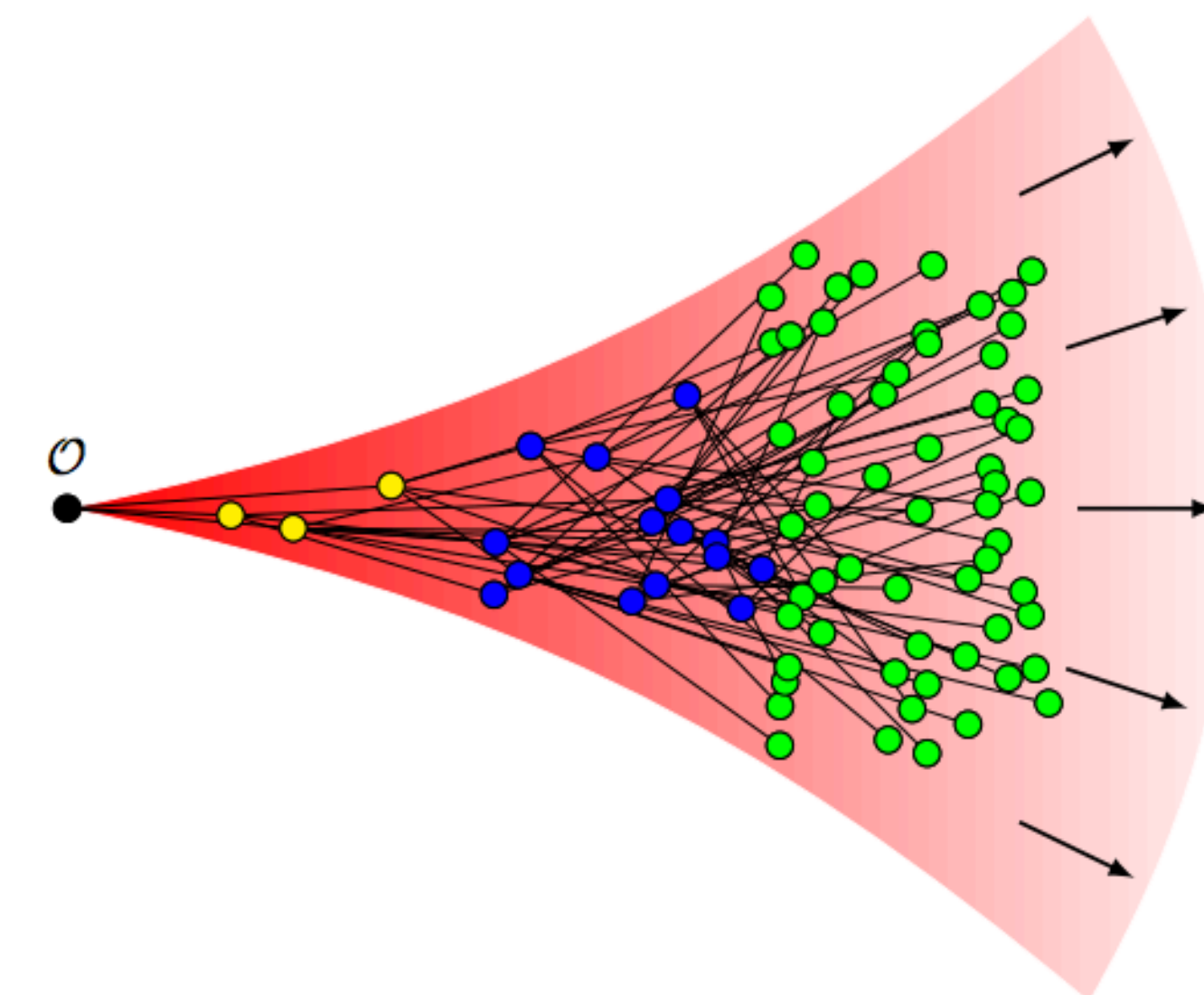
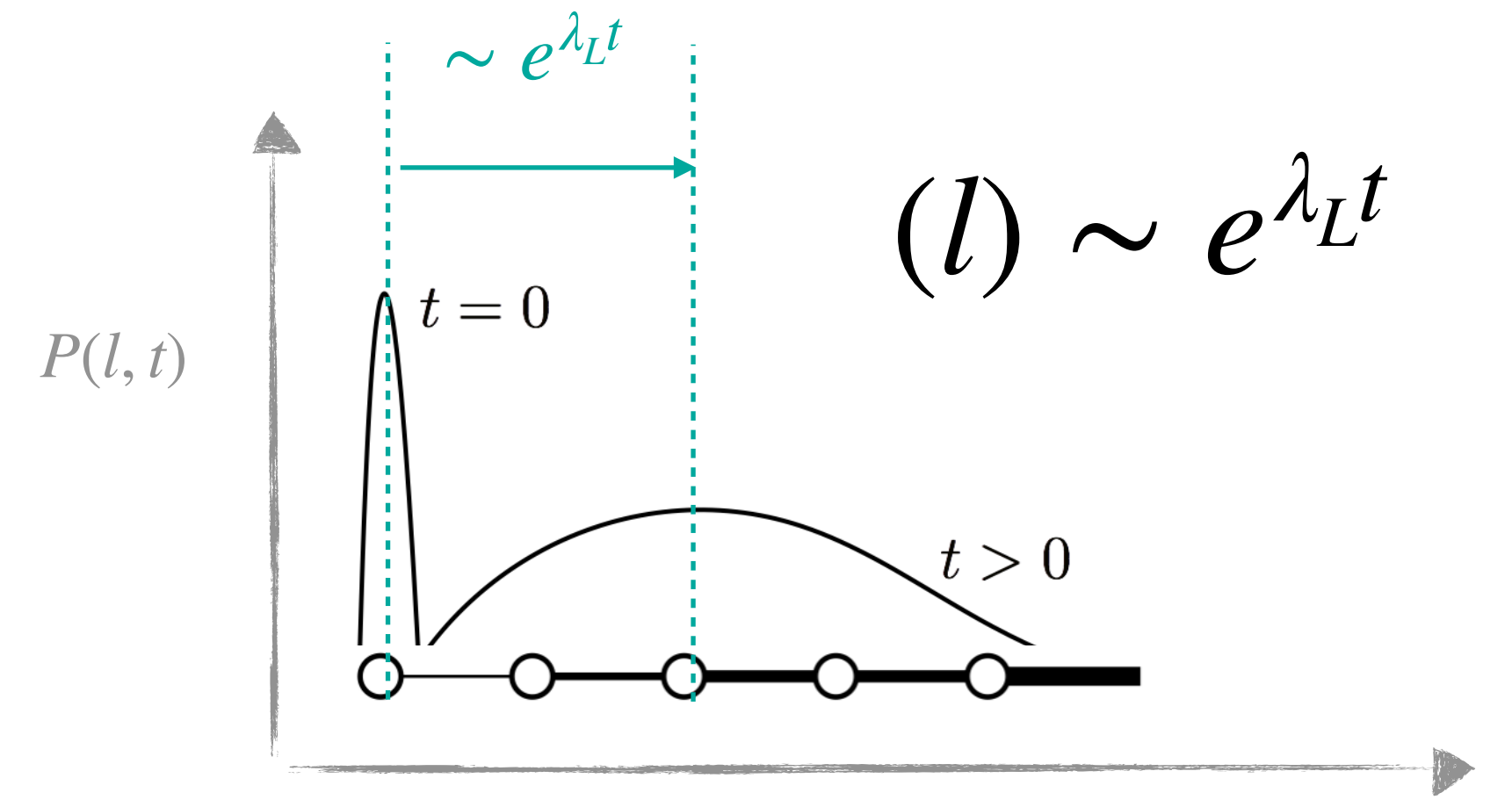


$$l \sim n^h$$

**Bound:**

$$h \equiv \lambda_L / 2\alpha \leq 1$$

**Size distribution**

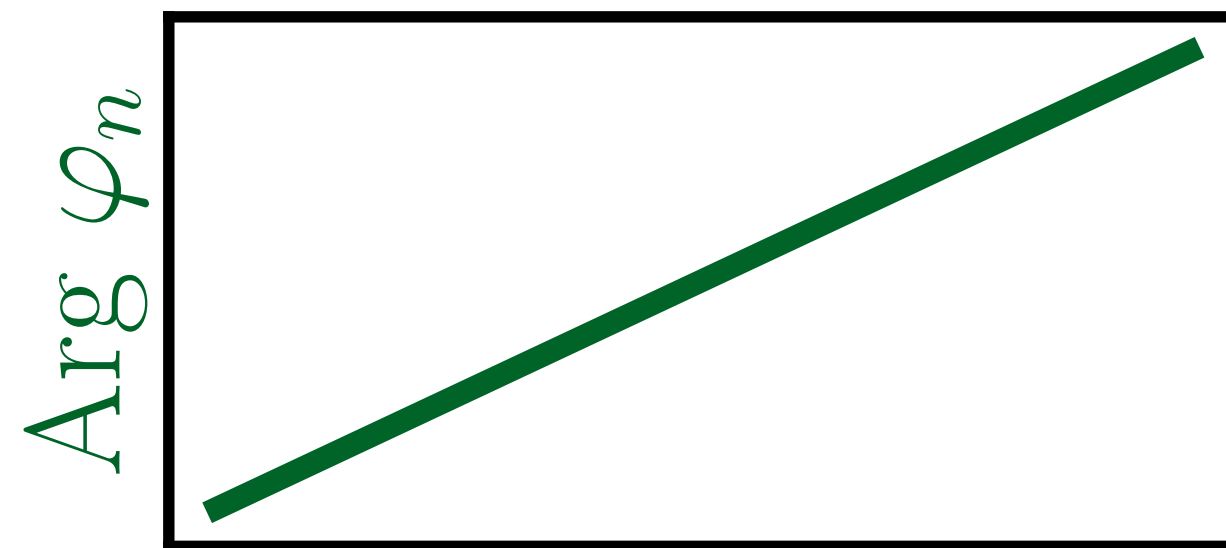


# Size Winding: Phase linearity

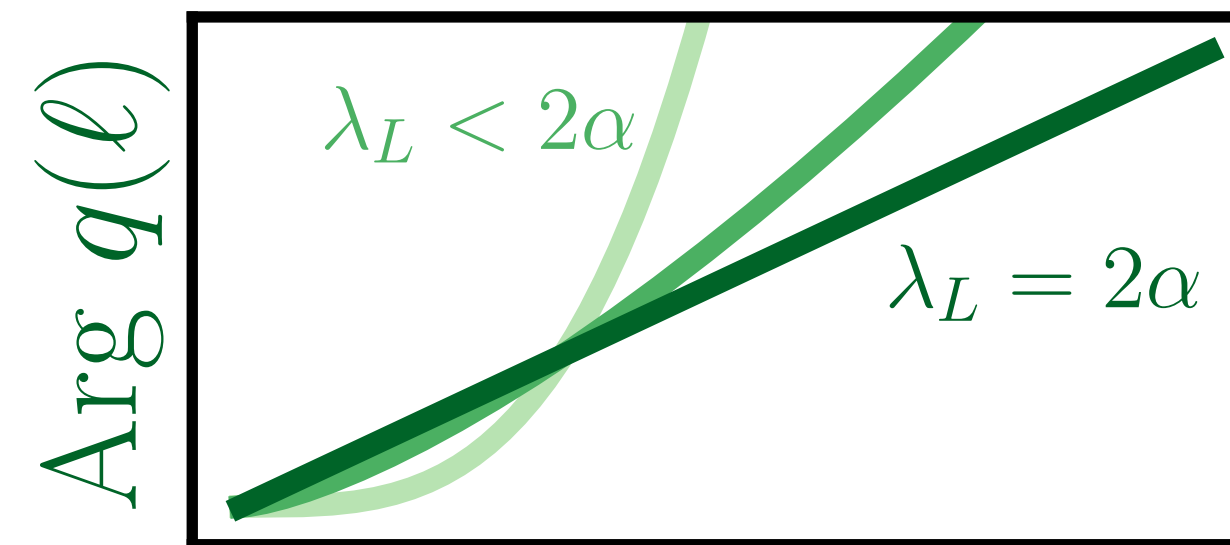
$$n \sim l^{1/h}$$

$$\text{Arg}[\varphi_n] = \theta(t)n \implies \text{Arg}[c_P] \sim \theta(t) l^{1/h}$$

$|P| = l$

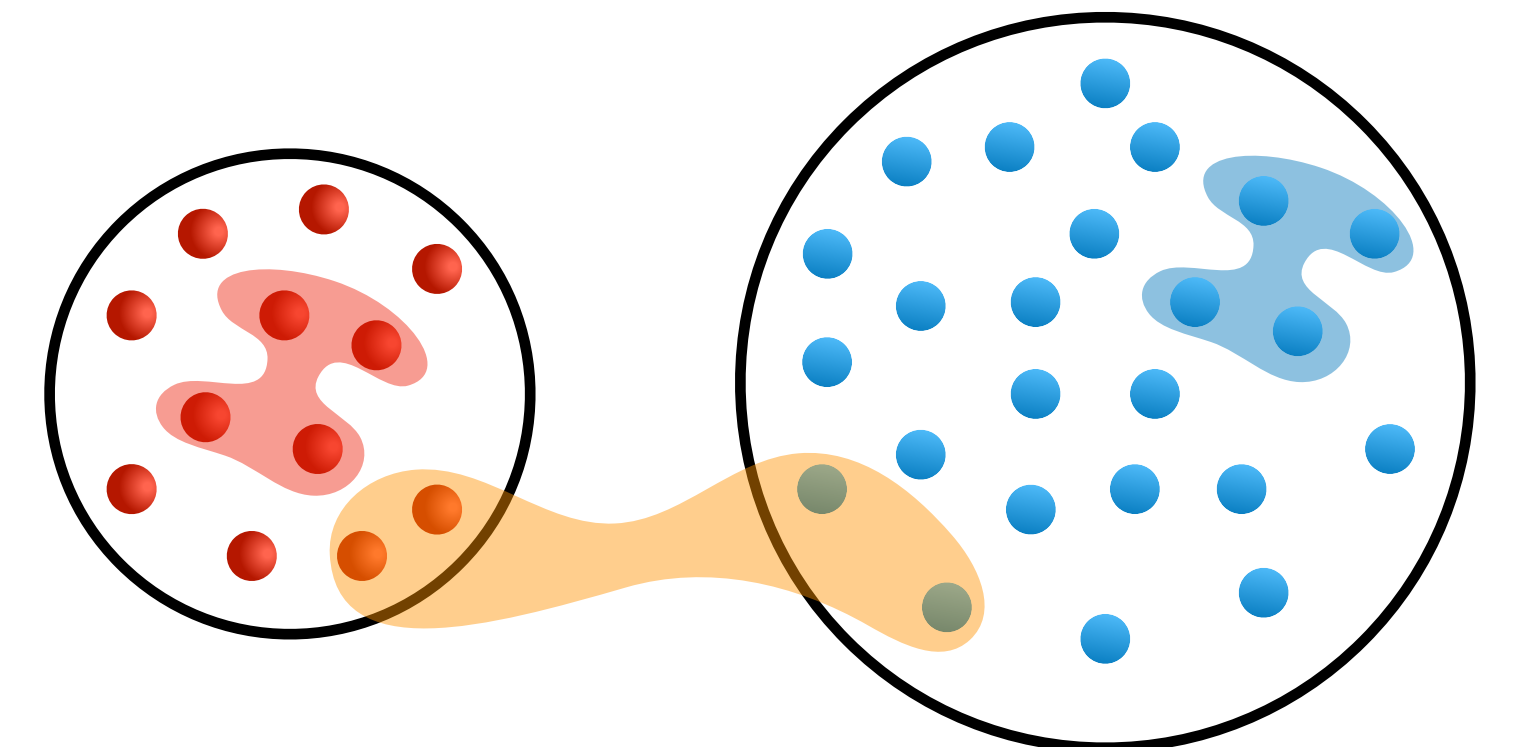


Krylov basis

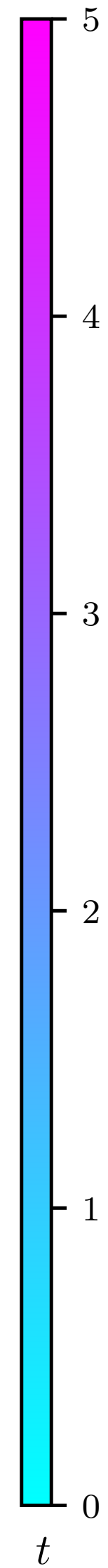
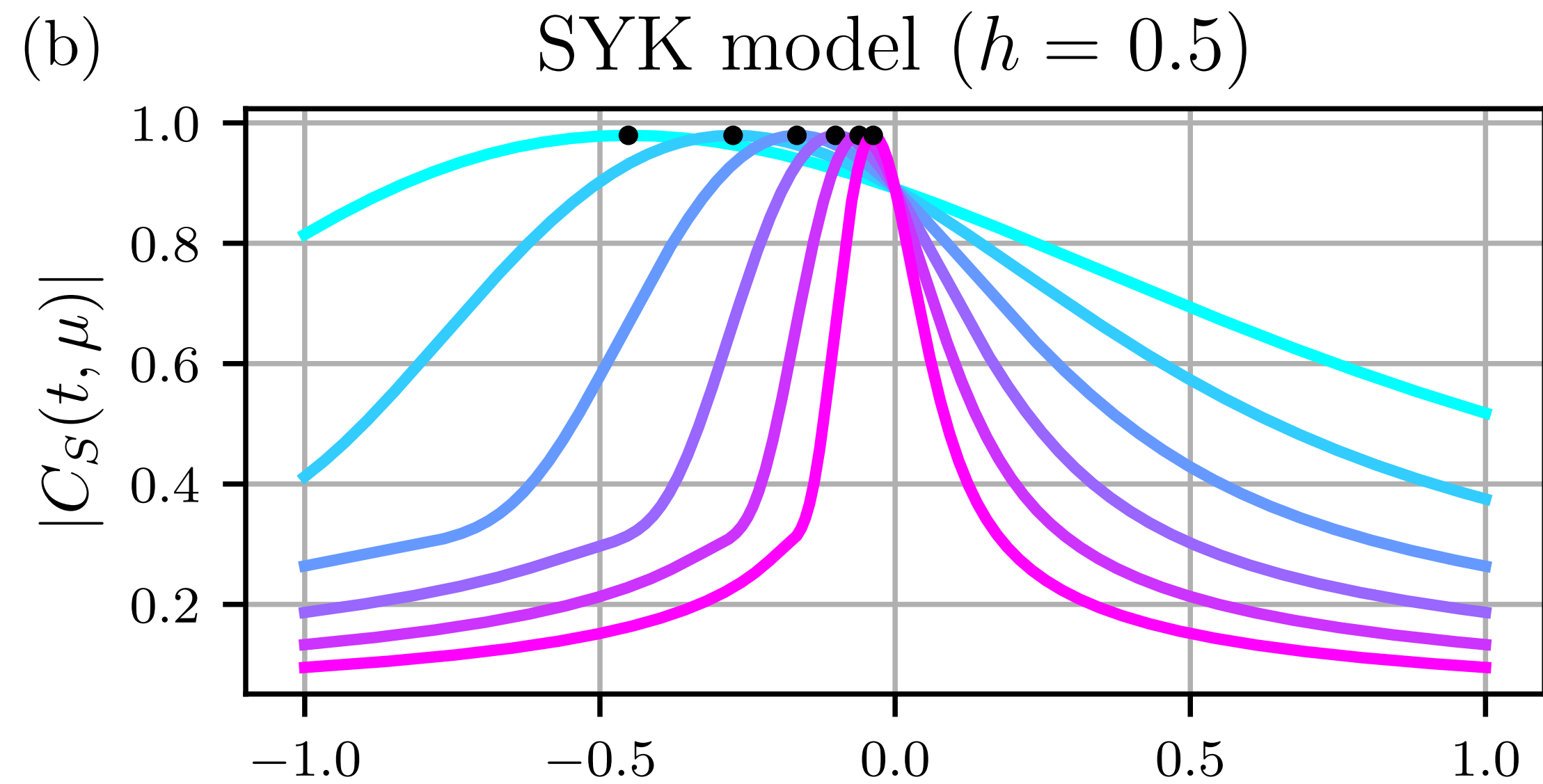
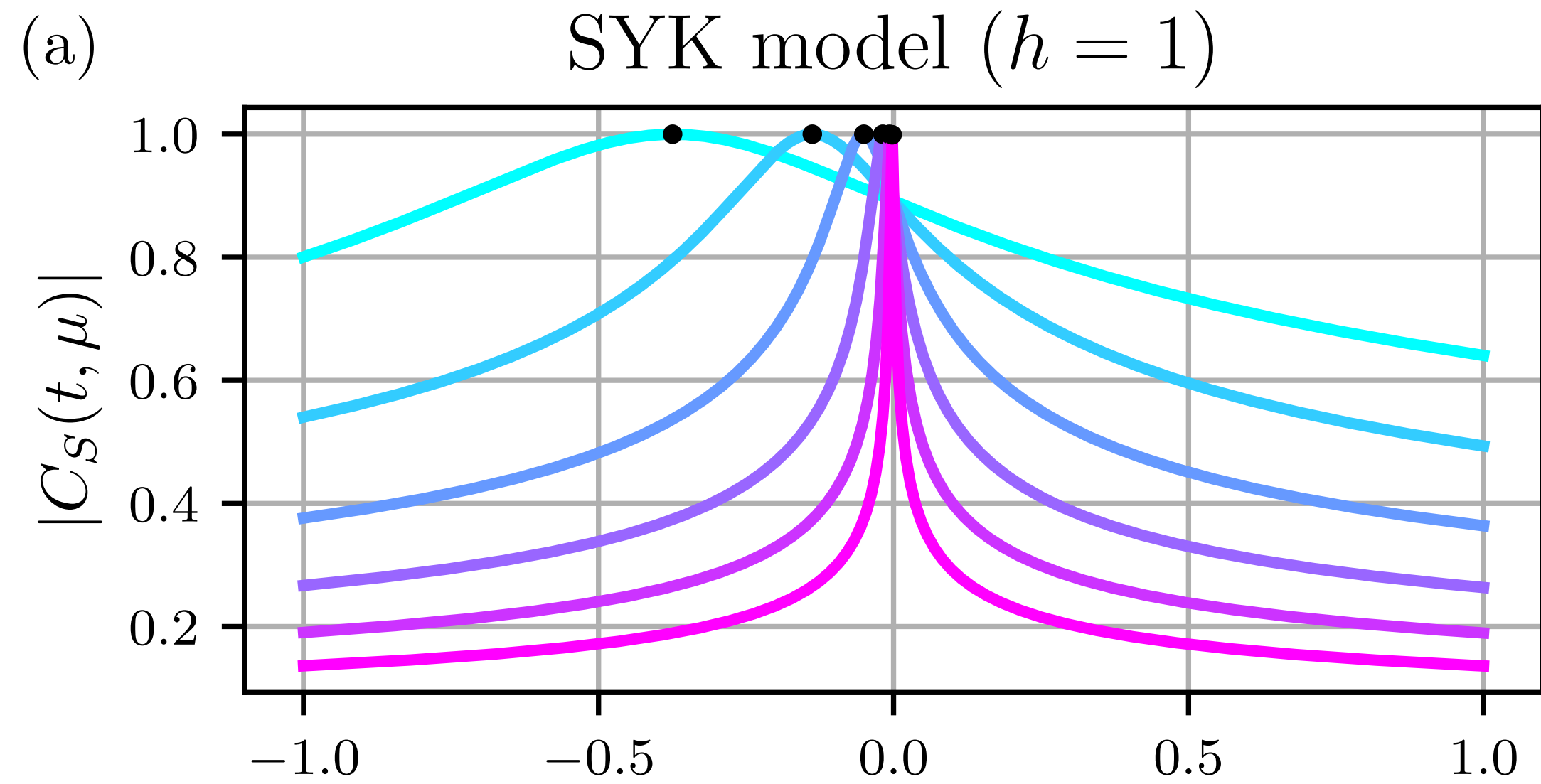


Size basis

This result is proven analytically in large- $q$  SYK model coupled to a bath



# Size Winding: SYK model



# Conclusion

- **Krylov winding is generic in non-integrable models:** it is a direct consequence of the operator growth hypothesis
- **Krylov winding  $\implies$  size winding** if two sufficient conditions are met
  - **Phase alignment** is guaranteed if the "size-resolved Krylov overlap matrix" is rank 1.
  - **Phase linearity** is guaranteed if the bound  $\lambda_L \leq 2\alpha$  (with  $\alpha$  the Krylov growth rate) is saturated. More generally,  $\text{Arg}[q(l)] \sim l^{1/h}$  with  $0 \leq h \equiv \lambda_L/2\alpha \leq 1$ .

# Thanks for your attention!

Davis Thuillier Omid Tavakol Thomas Scaffidi



Scaffidi group, APS Global Physics Summit 2025

Extra slides

# Size Winding: Phase alignment

“Size-resolved Krylov overlap matrix”:

$$M_{nm}(l) \equiv (O_n | \hat{P}_l | O_m)$$

with  $\hat{P}_l$  the projector onto the size- $l$  sector

[Chen, Mu, Wang, Zhang PRL 2025]

Size and winding distribution

$$P(l) = \sum_{P:|P|=l} |c_P(t)|^2 = \sum_{nm} \varphi_n^*(t) \varphi_m(t) M_{nm}(l)$$

$$Q(l) = \sum_{P:|P|=l} c_P(t)^2 = \sum_{nm} \varphi_n(t) \varphi_m(t) M_{nm}(l)$$

In large- $q$  SYK+bath,  $M_{nm}(l)$  has rank-1:

$$M_{nm}(l) = \psi_n(l) \psi_m(l)$$

[Chen, Mu, Wang, Zhang PRL 2025]

Perfect phase alignment:

$$\frac{|Q(l, t)|}{P(l, t)} = 1$$

More generally, (imperfect) phase alignment is guaranteed if  $\text{rank}[M_{nm}] = \simeq 1$

(i.e. if  $K$  has a single dominant eigenvalue).

Conjecture: this is true generically for  $k$ -local models.

# The Liouvillian graph

$$O(t) = e^{i\mathcal{L}t}O$$

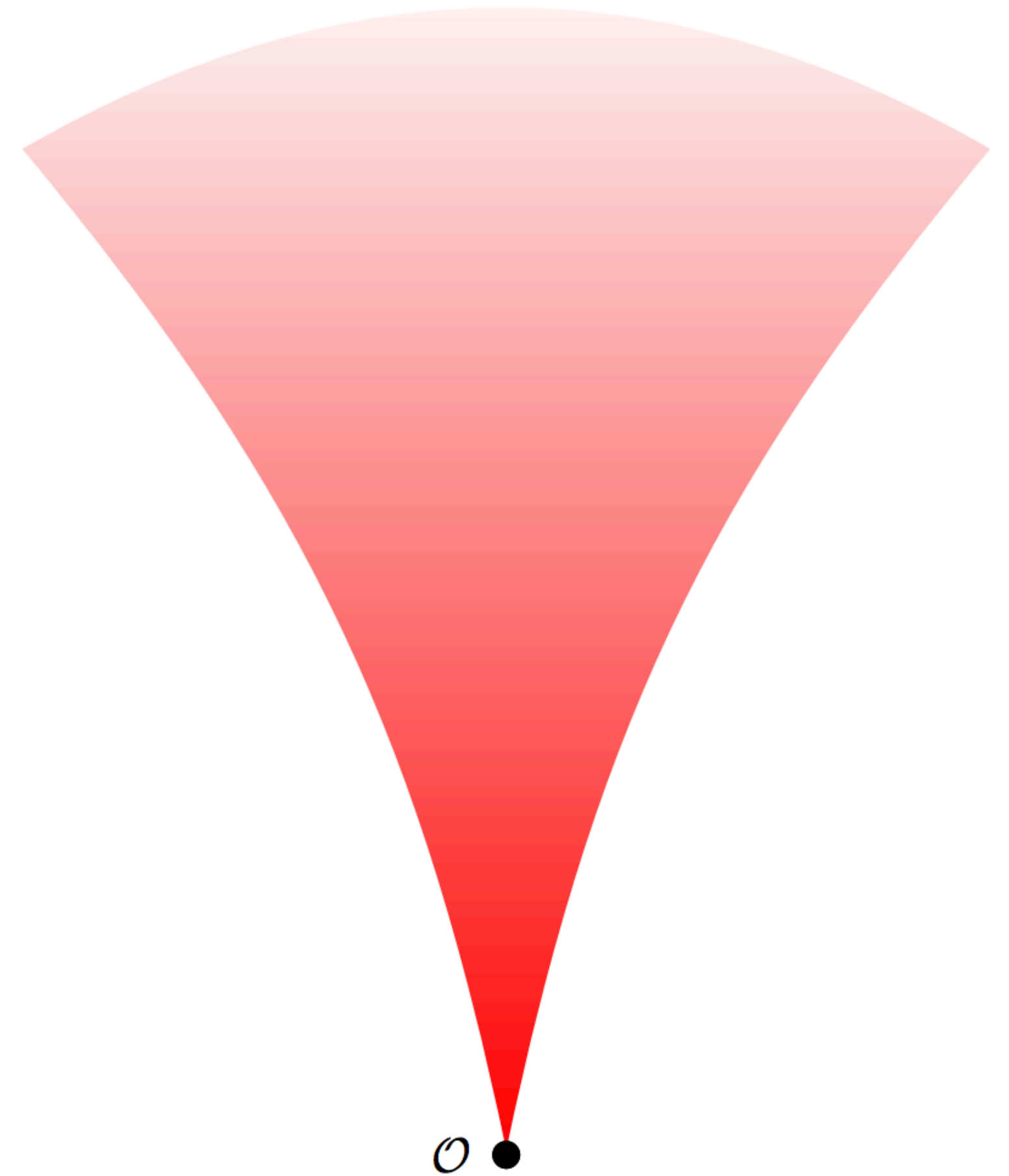
$$= O + (it)\mathcal{L}O + \frac{(it)^2}{2}\mathcal{L}^2O + \dots$$

$$\mathcal{L} = [H, \cdot]$$

**Example.** 1D Chaotic Ising model

$$H = \sum_i X_i + 1.05Z_iZ_{i+1} + 0.5Z_i$$

$$O = X_1$$



# The Liouvillian graph

$$\begin{aligned} O(t) &= e^{i\mathcal{L}t} \\ &= O + (it)\mathcal{L}O + \frac{(it)^2}{2}\mathcal{L}^2O + \dots \end{aligned}$$

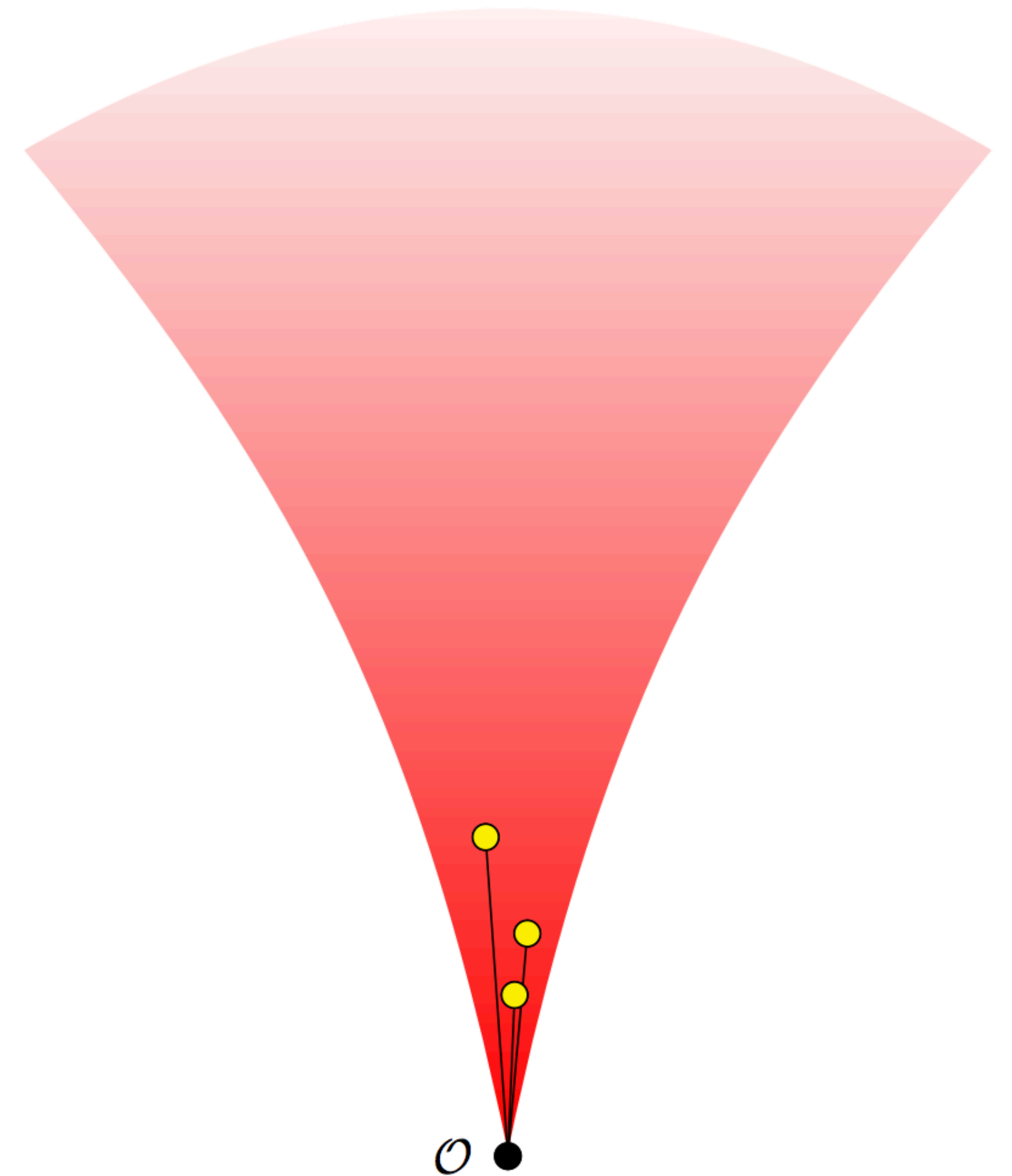
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**Example.** 1D Chaotic Ising model

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$$O = X_1$$

$$\mathcal{L}O = 1.05iY_1Z_2 + 1.05iZ_1Y_2 + 0.5iY_1$$



# The Liouvillian graph

$$\begin{aligned}
 O(t) &= e^{i\mathcal{L}t} \\
 &= O + (it)\mathcal{L}O + \frac{(it)^2}{2}\mathcal{L}^2O + \dots
 \end{aligned}$$

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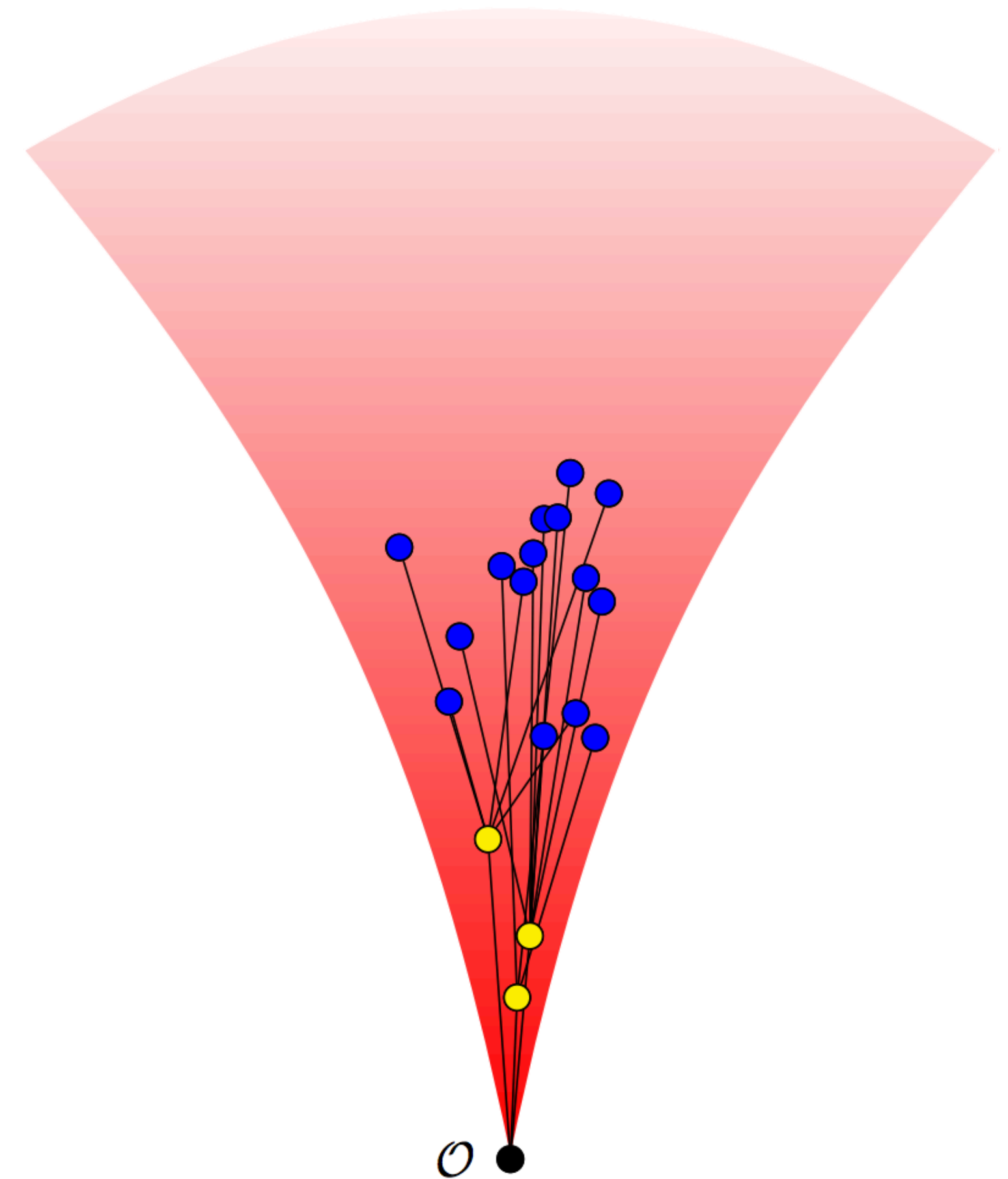
**Example.** 1D Chaotic Ising model

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$$O = X_1$$

$$\mathcal{L}O = 1.05iY_1Z_2 + 1.05iZ_1Y_2 + 0.5iY_1$$

$$\begin{aligned}
 \mathcal{L}^2O &= 2.1Z_1Z_2 - 2.1Y_1Y_2 + 0.25X_1 \\
 &+ 1.05^2Z_0X_1Z_2 + 1.05^2X_1 + 1.05^2X_2 \\
 &+ 1.05^2Z_1X_2Z_3 + 0.525X_1Z_2 + 0.525Z_1X_2
 \end{aligned}$$



# The Liouvillian graph

$$O(t) = e^{i\mathcal{L}t}$$

$$\mathcal{L} = [H, \cdot]$$

$$= O + (it)\mathcal{L}O + \frac{(it)^2}{2}\mathcal{L}^2O + \dots$$

**Example.** 1D Chaotic Ising model

$$H = \sum_i X_i + 1.05Z_iZ_{i+1} + 0.5Z_i$$

$$O = X_1$$

$$\mathcal{L}O = 1.05iY_1Z_2 + 1.05iZ_1Y_2 + 0.5iY_1$$

$$\begin{aligned} \mathcal{L}^2O &= 2.1Z_1Z_2 - 2.1Y_1Y_2 + 0.25X_1 \\ &+ 1.05^2Z_0X_1Z_2 + 1.05^2X_1 + 1.05^2X_2 \\ &+ 1.05^2Z_1X_2Z_3 + 0.525X_1Z_2 + 0.525Z_1X_2 \end{aligned}$$

